



# Stimulated Raman and Brillouin scattering

Sergey A. Ponomarenko

Dalhousie University



Back

Close

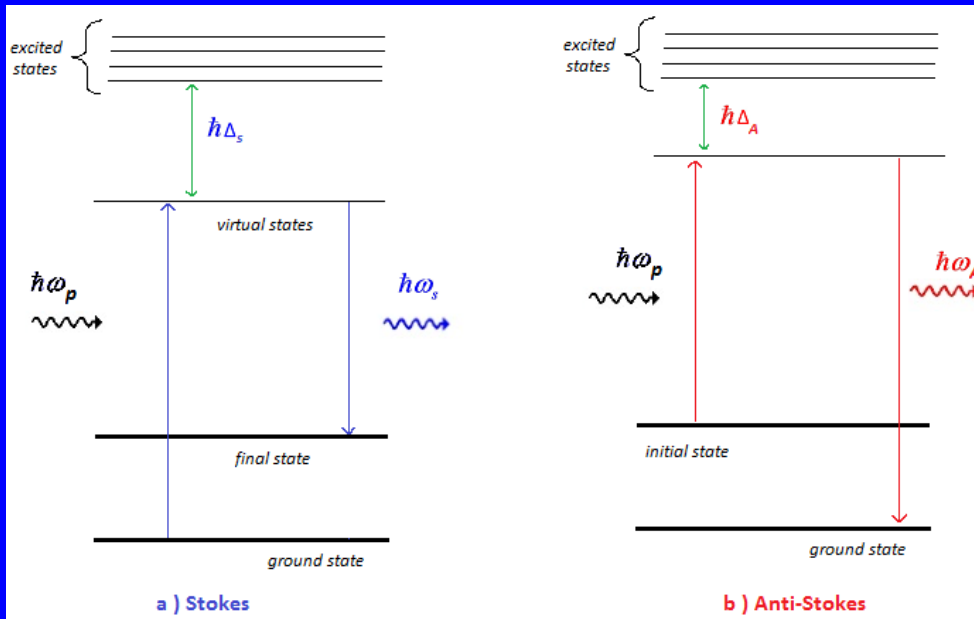


# Outline

- Basics of spontaneous Raman scattering
- Stimulated Raman scattering: cw case
- Spontaneous Brillouin scattering
- Stimulated Brillouin scattering (SBS): cw case
- Transient SBS
- SBS gain and threshold

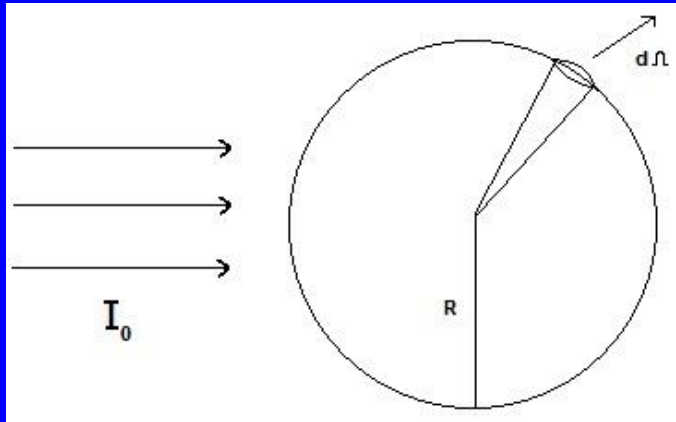


# Basics of Raman scattering



$$w_e \propto e^{-(E_f - E_g)/kT} \implies \text{anti-Stokes suppressed}$$





## Differential scattering cross-section (DSCS)

$$P = \sigma I_0$$

$$\frac{dP}{dS} = \frac{1}{R^2} \left( \frac{dP}{d\Omega} \right) = \frac{I_0}{R^2} \times \underbrace{\left( \frac{d\sigma}{d\Omega} \right)}_{\text{DSCS}}$$



No Stokes mode  $\implies$  spontaneous scattering

$$P_{\Delta} = \int_{d\Omega} \frac{dP}{d\Omega} \simeq NI_0 \Delta\Omega \frac{d\sigma}{d\Omega}$$

DSCS tabulated relative to nitrogen,

$$\left. \frac{d\sigma}{d\Omega} \right|_{N_2} = 5.5 \times 10^{-31} \text{ cm}^2/\text{Sr}/\text{molecule}$$

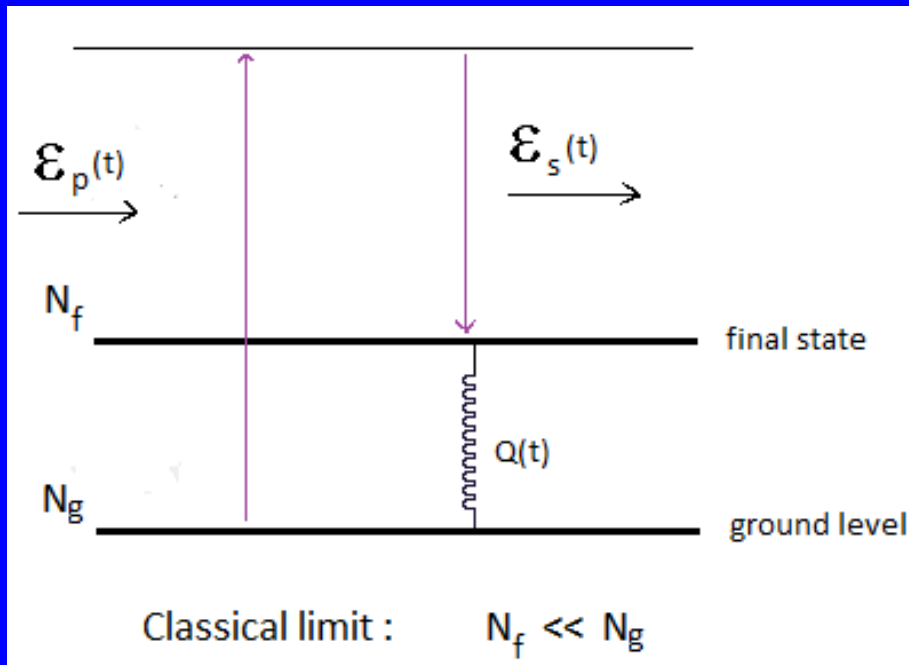
Conclusion:  $P \propto I_0$ , weak process!



# Classical theory of stimulated Raman scattering (SRS)



6/29



Back

Close

Potential energy of a polarized molecule:

$$W = -\frac{1}{2}pE$$

Dipole moment:

$$p = \epsilon_0 \underbrace{\alpha(\mathcal{Q})}_{\text{polarizability}} E$$

Polarizability:

$$\alpha \simeq \alpha_0 + \left( \frac{d\alpha}{d\mathcal{Q}} \right)_0 \mathcal{Q}$$

Force on a molecule:

$$F_d = -\nabla_{\mathcal{Q}} W$$





Harmonic equation of motion:

$$\partial_t^2 \mathcal{Q} + 2\gamma \partial_t \mathcal{Q} + \omega_0^2 \mathcal{Q} = F_d/m$$

SRS driving electromagnetic field (pump+Stokes):

$$E(t, z) = \frac{1}{2} [\mathcal{E}_p e^{i(k_p z - \omega_p t)} + \mathcal{E}_s e^{i(k_s z - \omega_s t)} + c. c.]$$

CW driven solution to molecular oscillator:

$$\mathcal{Q}(t, z) = \frac{1}{2} \mathcal{Q}_\omega e^{i(k_p - k_s)z} e^{-i\omega_\Delta t} + c. c$$

$$\mathcal{Q}_\omega = \frac{\epsilon_0 (d\alpha/d\mathcal{Q})_0 \mathcal{E}_p \mathcal{E}_s^*}{2m(\omega_0^2 - \omega_\Delta^2 - 2i\gamma\omega_\Delta)}$$

$\omega_\Delta = \omega_p - \omega_s$  Raman frequency shift



Back

Close





Induced polarization:

$$P_{NL} = N p_{NL} = \epsilon_0 N \left( \frac{d\mathcal{Q}}{d\alpha} \right)_0 \mathcal{Q} E$$

Pump and Stokes polarization components:

$$P_{NL} = \frac{1}{2} [\mathcal{P}_{NL}(\omega_s) e^{-i\omega_s t} + \mathcal{P}_{NL}(\omega_p) e^{-i\omega_p t} + c. c.]$$

Coupled-mode equations:

$$2ik_j \partial_z \mathcal{E}_j = -\mu_0 \omega_j^2 \mathcal{P}_{NL}(\omega_j) e^{-ik_j z}, \quad j = p, s$$

$$d_z I_p = -g_R I_p I_s$$

$$d_z I_s = g_R I_p I_s$$



Back

Close



Back

Close

Raman gain coefficient:

$$g_R = g_{R0} \mathcal{L}(\omega_\Delta)$$
$$g_{R0} = \frac{\omega_s N (d\alpha/d\mathcal{Q})_0^2}{8m\omega_0 \gamma c^2 n_p n_s}$$

Raman gain profile:

$$\mathcal{L}(\omega_\Delta) = \frac{\gamma^2}{[(\omega_\Delta - \omega_0)^2 + \gamma^2]}$$

Typical magnitude of center-line gain:  $g_{R0} \simeq 1.5$   
cm/GW, molecular hydrogen.



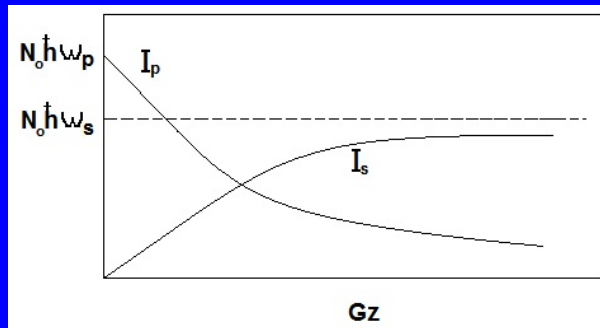


$$d_z I_p = -g_R I_p I_s$$

$$d_z I_s = g_R I_p I_s$$

Undepleted pump,  $I_p = \text{const}$ , exponential gain of Stokes mode!

Beyond undepleted pump approximation:



Back

Close

Driving anti-Stokes mode:

$$E(t, z) = \frac{1}{2} [\mathcal{E}_p e^{i(k_p z - \omega_p t)} + \mathcal{E}_s e^{i(k_s z - \omega_s t)} + \mathcal{E}_{as} e^{i(k_{as} z - \omega_{as} t)} + c.c.]$$

Stokes/anti-Stokes shifts:

$$\omega_{\Delta} = \omega_p - \omega_s = \omega_{as} - \omega_p$$

Material oscillator displacement:

$$\mathcal{Q}(t, z) = \frac{1}{2} \{ [\mathcal{Q}_{\omega_s} e^{i(k_p - k_s)z} + \mathcal{Q}_{\omega_{as}} e^{i(k_{as} - k_p)z}] e^{-i\omega_{\Delta} t} + c.c. \}$$

Nonlinear polarization:

$$P_{NL} = \frac{1}{2} [\mathcal{P}_{NL}(\omega_s) e^{-i\omega_s t} + \mathcal{P}_{NL}(\omega_{as}) e^{-i\omega_{as} t} + \mathcal{P}_{NL}(\omega_p) e^{-i\omega_p t} + c.c.]$$





Coupled-mode equations:

$$i\partial_z \mathcal{E}_s = -\xi_s (|\mathcal{E}_p|^2 \mathcal{E}_s + \mathcal{E}_p^2 \mathcal{E}_{as}^* e^{i\Delta k z})$$

$$i\partial_z \mathcal{E}_{as} = -\xi_{as} (|\mathcal{E}_p|^2 \mathcal{E}_{as} + \mathcal{E}_p^2 \mathcal{E}_s^* e^{i\Delta k z})$$

Phase mismatch:

$$\Delta k = 2k_p - k_s - k_{as}$$

Undepleted pump approximation:

$$I_p = \text{const}$$



Back

Close

Undepleted pump + large phase mismatch  $\implies$   
uncoupled Stokes/anti-Stokes modes:

$$d_z I_s = g_s I_p I_s$$

$$d_z I_{as} = -g_{as} I_p I_{as}$$

Exponential gain for Stokes and loss for  
anti-Stokes!





# Spontaneous Brillouin scattering

Light scattering by medium density fluctuations  
(sound waves):

$$\varepsilon' = \left( \frac{\partial \varepsilon}{\partial \rho} \right) \rho' = \frac{\gamma_e}{\rho} \rho'$$

electrostriction coefficient:

$$\gamma_e \equiv \rho \left( \frac{\partial \varepsilon}{\partial \rho} \right)$$

Induced polarization:

$$P_{NL} = \varepsilon_0 \varepsilon' E$$







Pump field:

$$E = \frac{1}{2} \mathcal{E} e^{i(\mathbf{k}_p \cdot \mathbf{r} - \omega_p t)} + c.c.$$

Sound wave:

$$\rho' = \frac{1}{2} \tilde{\rho} e^{i(\mathbf{q} \cdot \mathbf{r} - \Omega t)} + c.c.$$

Dispersion laws:

$$\omega_p = k_p c \quad \text{photons}$$

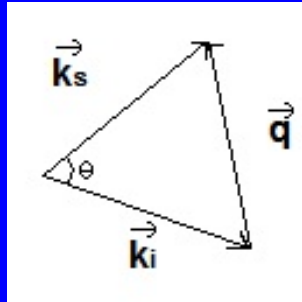
$$\Omega = q c_s \quad \text{phonons}$$

$c_s$  is sound speed.



Back

Close



$$\omega_s = \omega_p - \Omega$$

$$\mathbf{k}_s = \mathbf{k}_p - \mathbf{q}$$

$$\Omega \ll \omega_p \implies k_p \approx k_s:$$

$$\Omega = 2k_p \sin \theta / 2$$

Maximum generated phonon energy in the backward direction,  $\theta = \pi$ !



Back

Close



# Brillouin phonon propagation

Sound waves in 1D geometry:

- Mass conservation;
- Momentum conservation (second Newton's law)

Postulating,

$$\partial_t \rho + \partial_z(\rho v) = 0$$

$$\partial_t v + v \partial_z v = -\frac{1}{\rho} \partial_z p_{eff} + \left( \zeta + \frac{4}{3} v \right) \partial_{zz}^2 v$$



Back

Close

Effective pressure including “negative” electrostriction term:

$$p_{\text{eff}} = p - \gamma_e \frac{\epsilon_0 E^2}{2}$$

Adiabatic sound speed:

$$c_s = \sqrt{\left( \frac{\partial p}{\partial \rho} \right)_s}$$

Sound wave equation with  $E = 0$ :

$$\partial_{tt}^2 \rho' - c_s^2 \partial_{zz}^2 \rho' - \underbrace{\Gamma \partial_{tzz}^3 \rho'}_{\text{loss}} = 0$$





Sound wave propagation:

$$\rho' \propto e^{i(qz - \Omega t)}$$

Phonon dispersion relation:

$$q \simeq \Omega/c_s + i\alpha/2$$

Sound damping rate:

$$\alpha = \Gamma q^2 / c_s$$

Estimating,  $\Omega/2\pi \simeq 2$  GHz,  $\alpha^{-1} \sim 5\mu\text{m} \implies$

- strongly damped ultrasound SBS phonons
- negligible phonon displacements in SBS

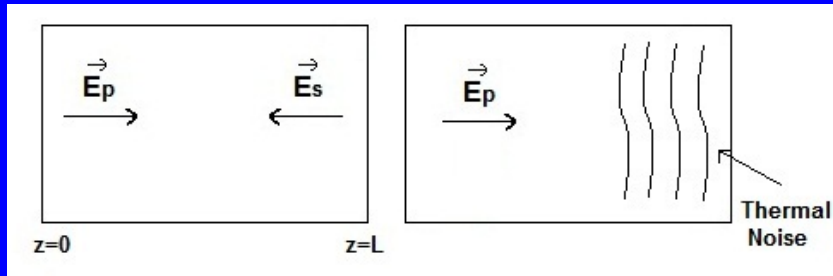


Back

Close



# Transient SBS (amplification mode)



Driving electromagnetic field:

$$E(t, z) = \frac{1}{2} [\mathcal{E}_p(t) e^{i(k_p z - \omega_p t)} + \mathcal{E}_s(t) e^{i(-k_s z - \omega_s t)} + c. c.]$$

Sound wave:

$$\rho' = \frac{1}{2} \tilde{\rho} e^{i(qz - \Omega t)} + c. c.$$



Back

Close

# Kinematics of photon-phonon interaction:

$$q = k_p + k_s \simeq 2k, \quad \Omega = \omega_p - \omega_s,$$

Density relaxation equation:

$$\partial_t \tilde{\rho} = -(\Gamma_B/2 + i\Delta) \tilde{\rho} + \left( \frac{i\varepsilon_0 \gamma_e k^2}{\Omega_B} \right) \mathcal{E}_p \mathcal{E}_s^*$$

Brillouin frequency detuning,

$$\Delta = \Omega_B - \Omega$$

$$\Omega_B = qc_s \quad \Gamma_B = \Gamma q^2$$



coupled-wave equations:

$$-\partial_z \mathcal{E}_s + \beta_s \partial_t \mathcal{E}_s = \left( \frac{i\omega\gamma_e}{4cn\rho_0} \right) \mathcal{E}_p \tilde{\rho}^*$$

$$\partial_z \mathcal{E}_p + \beta_p \partial_t \mathcal{E}_p = \left( \frac{i\omega\gamma_e}{4cn\rho_0} \right) \mathcal{E}_s \tilde{\rho}$$

coupled to the SBS phonon relaxation equation:

$$\partial_t \tilde{\rho} = -(\Gamma_B/2 + i\Delta) \tilde{\rho} + \left( \frac{i\varepsilon_0\gamma_e k^2}{\Omega_B} \right) \mathcal{E}_p \mathcal{E}_s^*$$

completely specify transient SBS!





## CW limit of SBS:

$$d_z I_s = -g_B I_s I_p,$$

$$d_z I_p = -g_B I_s I_p$$

$$g_B = g_{B0} \frac{\Gamma_B^2/4}{\Gamma_B^2/4 + \Delta^2}$$

$$g_{B0} = \frac{\omega^2 \gamma_e^2}{nc^3 c_s \rho_0 \Gamma_B}$$

center-line gain estimate  $g_{B0} \simeq 0.2$  cm/MW for  $\text{CS}_2$ , for example.

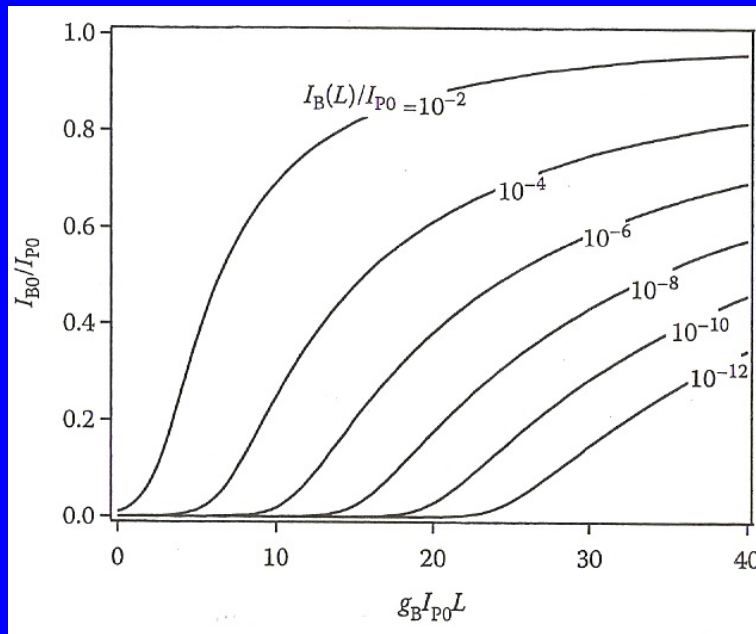




Undepleted pump  $I_p = \text{const}$ :

$$I_s(z) = I_s(L) \exp[g_B I_p (L - z)]$$

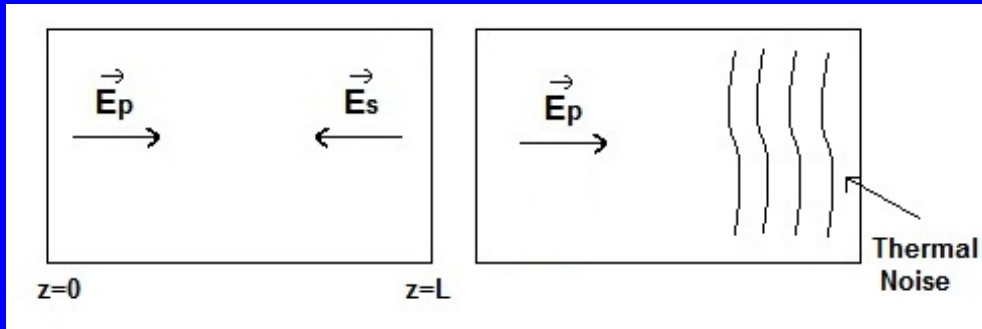
Beyond undepleted pump:



Back

Close

# SBS threshold



Metric: Reflection coefficient,

$$R = \frac{I_{s0}}{I_{p0}}$$

Stokes intensity at the exit:

$$I_{s0} = I_s(L)e^G$$



Overall gain from noise:

$$G = g_{B0} I_{p0} L$$

Universal gain interval:

$$25 \leq G_{th} \leq 30$$



# SBS threshold for cw laser sources

Parameters to consider: source power  $P$ , pump beam spot-size  $w_0$ ,

$$I_{p0} \simeq P / \pi w_0^2$$

Interaction length:

$$L \sim z_R \simeq k_p w_0^2$$

At threshold,

$$G = g_{B0} I_{p0} L = G_{th} \implies P_{th} \simeq \frac{G_{th} \lambda_p}{2g_{B0}}$$

Estimate:  $P_{th} \simeq 7$  kW for  $\text{CS}_2$ .

